

Geometric dualities and model theory

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Tarskian duality

Theory $T \longleftrightarrow$ Class of models $\mathfrak{M}(T)$

For a κ -categorical T

Theory $T \longleftrightarrow$ Model M_T (of cardinality κ)

Geometric dualities

Affine commutative \mathbb{C} -algebra

Complex algebraic variety

$$R = \mathbb{C}[X_1, \dots, X_n]/I$$

$$\mathbf{V}_R$$

Commutative unital C^* -algebra

Compact topological space

$$A$$

$$\mathbf{V}_A$$

Affine k -scheme

The geometry of k -definable points, curves etc of an algebraic variety \mathbf{V}_R

$$R = k[X_1, \dots, X_n]/I$$

k -scheme of finite type

The geometry of k -definable points, curves etc of a “Zariski geometry” \mathbf{V}_S

$$S$$



From the scheme – language – theory to the structure



Claim A

These are **syntax – semantics** dualities.

The dualities can be recast in the form of Tarskian dualities.

In general the syntax may come with a topology (as in C^* -algebras).

Recall also the syntax of *continuous model theory*.

What is geometric semantics?

Non-example. The models of the theory of arithmetic **do not** provide a semantics of geometric type.

Suggestions. The structures on the right hand side of dualities must be **stable** (in a generalised sense).

There also should be a topology associated in a natural way with the syntax.

This leads again to the notion of Zariski geometry.

L -Zariski geometries.

\mathbf{V} is said to be a Noetherian L -Zariski if it satisfies

- Closed subsets of V^n are exactly those which are L -positive-quantifier-free definable.
(There may be points which are not closed!)
- The projection of a closed set is constructible (positive quantifier-elimination).
- A good dimension notion on closed subsets is given. The addition formula and the fibre conditions hold.
- For every extension $L(C)$ of the language by constants, any $\mathbf{V}' \succeq \mathbf{V}$ as an $L(C)$ -structure satisfies the above.

Theorem. Noetherian L -Zariski geometries are of finite Morley rank.

Example. Affine k -schemes.

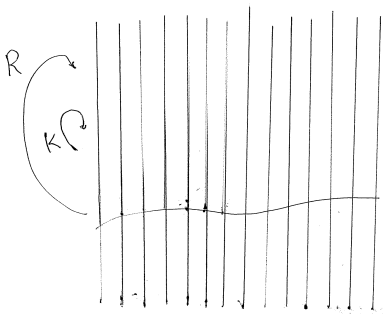
Given a field k and an affine k -algebra R we introduce a language L_R and an L_R -Zariski structure \mathbf{V}_R .

The language L_R of two-sorted structures:

- sort K for an algebraically closed field K containing k with names for elements of k ;
- sort V fibred as the union of all irreducible (1-dimensional) representations of R ;
- L_R has means to describe the action of the field K and the additive structure on fibres of V ;
- L_R has a name a for each $a \in R$ interpreted as a linear operator on each fibre (each vector space) of V .

The structure \mathbf{V}_R has exactly one fibre for each isomorphism type of irreducible representations.

Illustration: sort V and a section



We define natural morphisms between the L_R -structures (for different R).

Duality Theorem. *The category of affine k -algebras for all $k \subseteq K = K^{alg}$ is isomorphic to the category of respective Zariski structures.*

Points (non necessarily closed) of \mathbf{V}_R correspond to irreducible K -representations of R .

A noncommutative duality Theorem

The above duality can be extended to non-commutative geometry “at roots of unity”.

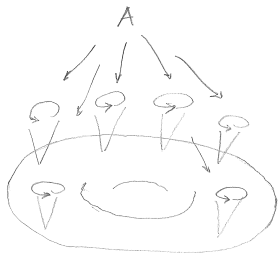
$$A_{\mathbf{V}} \longleftrightarrow \mathbf{V}_A.$$

$A_{\mathbf{V}}$ – co-ordinate algebra, \mathbf{V}_A – Zariski geometry.

A non-commutative example “at root of unity”

Non-commutative 2-torus at $q = e^{2\pi i \frac{m}{N}}$ has
 co-ordinate ring $A = A_q =$
 $\langle U, V : U^* = U^{-1}, V^* = V^{-1}, UV = qVU \rangle$

Points have structure of (an orthonormal basis)
 of a N -dim Hilbert space.



Affine commutative \mathbb{C} -algebra R

Complex algebraic variety \mathbf{V}_R

Commutative C^* -algebra A

Compact topological space \mathbf{V}_A

Affine k -scheme R

The k -definable structure on an algebraic variety \mathbf{V}_R

k -scheme of finite type S

The k -definable structure on a Zariski geometry \mathbf{V}_S

C^* -algebra A at roots of unity

Zariski geometry \mathbf{V}_A

Weyl-Heisenberg algebra

$\langle Q, P : QP - PQ = i\hbar \rangle$

?

The integers \mathbb{Z}

?

$$QP - PQ = i\hbar$$

“The whole of quantum mechanics is in this **canonical commutation relation**”.

An analogy:

$$Y = X^2 + aX + b$$

is (the equation of) a parabola.

$$H = \frac{1}{2}(P^2 + \omega^2 Q^2)$$

is (the Hamiltonian of) a quantum harmonic oscillator.

$$QP - PQ = i\hbar$$

This does not allow the C^* -algebra (Banach algebra) setting. On suggestion of Herman Weyl and following Stone – von Neumann Theorem replace the Weyl-Heisenberg algebra by the category of **Weyl $*$ -algebras**

$$A_{a,b} = \left\langle U^a, V^b : U^a V^b = e^{iab\hbar} V^b U^a \right\rangle,$$

$$a, b \in \mathbb{R}, U^a = e^{iaQ}, V^b = e^{ibP}.$$

where it is also assumed that U^a and V^b are unitary.

We may assume that $\frac{\hbar}{2\pi} \in \mathbb{Q}$ and so, when $a, b \in \mathbb{Q}$ the algebra $A_{a,b}$ is **at root of unity**. We call such algebras **rational Weyl algebras**.

Sheaf of Zariski geometries over the category of rational Weyl algebras

The category \mathcal{A}_{fin} has objects $A_{a,b}$, rational Weyl algebras, and morphisms = embeddings.

Note,

$$A_{a,b} \subseteq A_{c,d} \text{ iff } cn = a \ \& \ dm = b \text{ for } n, m \in \mathbb{Z} \setminus \{0\}.$$

This corresponds to the surjective morphism in the dual category \mathcal{V}_{fin} of Zariski geometries

$$\mathbf{V}_{A_{a,b}} \rightarrow \mathbf{V}_{A_{c,d}}.$$

The duality functor

$$A \mapsto \mathbf{V}_A$$

can be interpreted as defining a **sheaf of Zariski geometries** over the category of rational Weyl algebras.

Completions of \mathcal{A}_{fin} and \mathcal{V}_{fin} .

The completion of \mathcal{A}_{fin} is \mathcal{A} , the category of all Weyl algebras in the Lie-groups topology (not the Banach algebra topology):

$$\sigma U^\alpha V^\beta \rightarrow sU^a V^b \text{ iff } \sigma \rightarrow s, \alpha \rightarrow a \text{ and } \beta \rightarrow b,$$

$$\sigma, \alpha, \beta \in \mathbb{Q}, s, a, b \in \mathbb{R}$$

Completing \mathcal{V}_{fin} is the main difficulty of the project.

We use **structural approximation**, which in basic cases is equivalent to the ultraproduct construction of **continuous model theory**.

Exercise. Use the ultraproduct of continuous model theory to construct the universal cover of the algebraic torus \mathbb{C}^\times .

The space of states.

We construct a \mathcal{V}_{fin} -**projective** limit object $\mathbf{V}_{\mathcal{A}}$, corresponding to the algebra $\bigcup \mathcal{A}$.

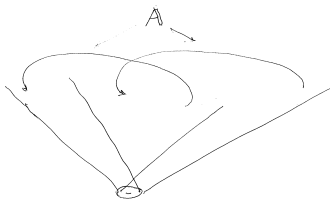
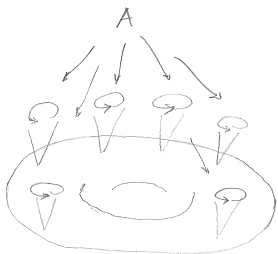
By construction

$$A \in \mathcal{A}_{\text{fin}} \Rightarrow \mathbf{V}_{\mathcal{A}} \twoheadrightarrow \mathbf{V}_A.$$

This object is closely related to *the Hilbert space of quantum mechanics*. We call the object **the space of states**.

Remark. The same construction for the category of commutative algebras $\langle U^a, U^{-a} \rangle$, $a \in \mathbb{Q}$, (equivalent to the category of étalé covers of \mathbb{C}^\times) produces the universal cover $(\mathbb{C}, +)$ of \mathbb{C}^\times and the morphism $\exp : \mathbb{C} \rightarrow \mathbb{C}^\times$.

How noncommutative \mathbf{V}_A deforms into $\mathbf{V}_{\mathcal{A}}$.



Operators P, Q on $\mathbf{V}_{\mathcal{A}}$.

We define in each member of the ultraproduct

$$Q := \frac{U^a - U^{-a}}{2ia}, \quad P := \frac{V^b - V^{-b}}{2ib}$$

in accordance with

$$U^a = e^{iaQ}, \quad V^b = e^{ibP}.$$

Then in the limit, we can calculate in $\mathbf{V}_{\mathcal{A}}$, for any state e

$$(QP - PQ)e = i\hbar e.$$

So, in the space of states:

$$QP - PQ = i\hbar I.$$

In other words, there is a Lie algebra $\langle Q, P \rangle$ acting on $\mathbf{V}_{\mathcal{A}}$ (the Heisenberg algebra)

Time evolution operators on the space of states

Theorem. *Automorphisms of \mathcal{A}_{fin} give rise to certain operators on $\mathbf{V}_{\mathcal{A}}$. These are definable in the sense of continuous model theory.*

Such operators $K (= K_{\text{particle}})$ are typically the “time evolution operators for a given particle”.

The one-parameter subgroups $\{K^t : t \in \mathbb{R}\}$ describe the time evolution of *the particle corresponding to K* .

Restriction to a commutative algebra case

Theorem. A “time evolution operator” on $(\mathbb{C}, +)$ (as the cover of \mathbb{C}^\times) has the form

$$z \mapsto \kappa z, \quad \text{some } \kappa \in \mathbb{R}_{>0}.$$

This corresponds to ‘raising to power’ κ on \mathbb{C}^\times .

Recall: the theory of raising to real power is superstable, provided Schanuel’s conjecture is true.

Example. Quantum harmonic oscillator.

The Hamiltonian:

$$H = \frac{1}{2}(P^2 + Q^2)$$

The time evolution operator :

$$K^t = K_H^t := e^{-i\frac{H}{\hbar}t}, \quad t \in \mathbb{R}.$$

This “induces” the automorphism of the category of algebras

$$U^a \mapsto e^{-\frac{2\pi a^2 \sin t \cos t}{2}} U^a \sin t V^a \cos t$$

$$V^a \mapsto e^{\frac{2\pi a^2 \sin t \cos t}{2}} U^{-a} \cos t V^a \sin t$$

Scheme of calculations

- rewrite the formula over $\mathbf{V}_{\mathcal{A}}$ in terms of Zariski-regular pseudo-finite sums and products over \mathbf{V}_A , $A \in \mathcal{A}_{fin}$;
- calculate **uniformly** in \mathbf{V}_A (using e.g. the *Gauss quadratic sums* formula)
- apply ultraproduct (continuous model theory) to the result and get the result in terms of the standard reals.

Example. Quantum harmonic oscillator.

$$K^t = K_H^t := e^{-i\frac{H}{\hbar}t}, \quad t \in \mathbb{R}.$$

To calculate K^t we *approximate* assuming $\sin t, \cos t \in \mathbb{Q}$. This transfers us to the rational category \mathcal{V}_{fin} and to the calculations in (pseudo)finite-dimensional spaces.

Then the matrix element on row x_1 and column x_2 (*kernel of the Feynman propagator*) is calculated as

$$\langle x_1 | K^t x_2 \rangle = \sqrt{\frac{1}{2\pi i \hbar \sin t}} \exp i \frac{(x_1^2 + x_2^2) \cos t - 2x_1 x_2}{2\hbar \sin t}.$$

The trace of K^t ,

$$\mathrm{Tr}(K^t) = \int_{\mathbb{R}} \langle x | K^t x \rangle = \frac{1}{\sin \frac{t}{2}}.$$

$$\mathrm{Tr}(K^t) = \int_{\mathbb{R}} \langle x | K^t x \rangle = \frac{1}{\sin \frac{t}{2}}.$$

Note that in terms of conventional mathematical physics we have calculated

$$\mathrm{Tr}(K^t) = \sum_{n=0}^{\infty} e^{-it(n+\frac{1}{2})},$$

a non-convergent infinite sum.

Work in progress

Given a Hamiltonian in the form

$$H_f = \frac{P^2}{2\hbar} + f(Q), \quad f(X) \in \mathbb{Q}[X]$$

one can calculate using the approximation method and setting $\Delta t = \frac{1}{\nu}$, ν a big non-standard integer,

$$K_f^{\Delta t} \approx e^{-i\frac{P^2}{2\hbar}\Delta t} \cdot e^{-if(Q)\Delta t}.$$

Both operators on the right **are definable**.

Given $t \in \mathbb{R}$, $t \approx \frac{\kappa}{\nu}$ it is natural to identify the time evolution operator for H_f on $\mathbf{V}_{\mathbb{A}}$ as

$$K_f^t := \left(e^{-i\frac{P^2}{2\hbar\nu}} e^{-i\frac{f(Q)}{\nu}} \right)^{\kappa}$$

(cf the Trotter product formula).

$$K_f^t := \left(e^{-i\frac{p^2}{2\hbar\nu}} e^{-i\frac{f(Q)}{\nu}} \right)^\kappa$$

Is it a regular transformation? I.e. is the image of the operator K_f^t (under standard part map) well-defined in the space of states?

Equivalently, does

$$\lim \left(e^{-i\frac{p^2}{2\hbar\nu}} e^{-i\frac{f(Q)}{\nu}} \right)^\kappa$$

exist?

This is equivalent to the Feynman path integral hypothesis for QM.